

TWO-DIMENSIONAL MATHEMATICAL MODEL OF FLOWS IN THIN FILM COMPOSITE MEMBRANES

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Abstract: Mathematical modelling of reverse osmosis membranes has evolved from simplified one-dimensional simulations to complex three-dimensional simulations using CFD based techniques. These models have been useful in simulating solute and solvent flows across the membrane, the development of the concentration polarisation layer and the effects of spacer and spacer geometry as some examples. Various simplifying assumptions are, however, made in the modelling process that limit their extension to the specific application of directly-coupled wave powered desalination. These include the treatment of the membrane and use of rejection coefficients for solute transport. The model presented in this paper addresses some of the limitations imposed currently on available models. Fully coupled mass-momentum equations are specified for the hydrodynamics within the feed, membrane and permeate channels. Semi-empirical relationships are developed to account for the effects of inlet pressure, cross-flow velocity and inlet concentration on solvent and solute flows across the membrane. The model is validated against published experimental data and the predicted errors for simulations of solvent and solute flows were found to be 0.6% and 0.7% respectively. The effects of three feed spacer types, submerged, cavity and zigzag, on solvent and solute flows are then considered. Larger wall shear stress was seen for the submerged type spacer than for either cavity or zigzag types.

Keywords: Computation Fluid Dynamics, Concentration Polarisation, Diffusion Coefficients, Reverse Osmosis, TFC membranes

https://doi.org/10.47412/GAEH3878

1. Introduction

Numerous mathematical models have been developed to characterise the reverse osmosis process [1-8]. These models can be broadly classified as either being analytical or numerical [9]. Independent of the classification, three areas must be considered when modelling the reverse osmosis process: (i) solute and solvent flows across the membrane; (ii) the effects of concentration polarisation and (iii) the hydrodynamics within the feed and permeate channels. The coupled nature of these areas has contributed to the complexity of the models that have been developed to accurately simulate the process. Although some models vary in their treatment of solvent and solute transport across the membrane, the governing equations are similar

[10]. The equations for solvent and solute flows posit a linear relationship between solvent and solute flows and their respective driving potentials, pressure and concentration, although in practice this is not the case [3, 4, 11-15]. Zhou & Song [14] presented experimental data on the increasing non-linear behaviours for both solvent and solute flows for increasing salt concentrations.

Considering solvent flows, Paul [13] noted that the linear relationship was untenable since the concentration gradient induced by the pressure difference must be bounded; a flaw that arises because of the subtle simplification implied in the derivation of solvent flow. Additionally, the assumption of a constant water permeability carried inaccuracies since it is affected by experimental variables such as pH, pressure and ionic strength [16]. In the instance of non-linearity of solute flows, limits to solute rejection have been noted in the literature [14, 17]. Assuming that the primary mechanism of solute transport through the active layer is diffusive, the diffusion coefficient is inherent to the diffusing material and can be affected by parameters such as viscosity, temperature and concentration [18-19]. Additionally, Countinho [20] noted that the increase in pressure resulted in a greater convection rate for the transport of solute to the membrane surface. This increases the concentration at the interface, causing an increase in the diffusivity of the solute in the opposite direction to that of the process pressure.

Concentration polarisation is a phenomenon where solute rejected by the membrane forms a layer of high concentration near the surface of the membrane, reducing the solvent flow across the membrane [21-23]. Various approaches have been proposed to estimate the concentration at the membrane surface [21, 23-27]. These models typically introduce a rejection coefficient to represent the selectivity of the membrane, opting to treat the membrane as a homogenous material. Most models do not take into account the various layers that comprise the membrane, though there have been some limited attempts at characterising the role played by the individual layers [17, 28]. There is a belief that the porous support structure may contribute significantly to the diffusive transport through the composite structure. The ratio of "real" (effective) flow and "ideal" flow is identical to the inverse ratio of the "ideal" and "real" diffusive path lengths [17]. This model however assumes that the porous support is comprised of uniform cylindrical pores, which in practice is not necessarily the case.

The hydrodynamics within the feed channel is affected by the feed spacer arrangement used. Feed spacers are typically composed of two sets of straight non-woven parallel cylindrical filaments [29]. Feed spacers form integral parts of reverse osmosis (RO) spiral wound modules for two purposes. Firstly, feed spacers keep membranes apart from the flow channel. Secondly, feed spacers promote mixing of the fluid, hence reducing the effects of concentration polarisation [30]. The use of computational fluid dynamics (CFD) based techniques has further enhanced our understanding of the hydrodynamics within the feed and permeate channels. Various authors have and continue to publish on the effects of feed spacers on solute and solvent flows through the membrane [31]. Broadly speaking, spacer models typically exclude permeation as this has little effect on wall shear stress [32], though examples are found in the literature [39-40]. Calculation of local permeate flows can increase computational load.

The mathematical model presented in this paper addresses the non-linear behaviour of solvent and solute flows across the membrane, concentration polarisation and the hydrodynamics within the feed and permeate channels. Fully coupled mass-momentum equations are used to define the hydrodynamics within the feed and permeate layers. Equations are specified for solute transport though the individual layers that comprise thin-film membrane. Further, consideration is given to the effects of inlet pressure, cross-flow velocity and concentration on solvent and solute flows across the membrane. Simulations were run for the effects of feed spacers on permeation through the membrane. It is noted that the work presented in this paper forms part of ongoing research into directly coupled wave-powered desalination [35-36]. The consideration of existing limitations is of significant importance to modelling this process.

2. Model Development

The basic model equations were previously presented in Maharajh & Persad [35]. As further modifications were made to the original model and for ease of reference, they are presented anew.

2.1 Fluid and Solute Flow in the Feed and Permeate Channels

The governing equations for fluid and solute flows within the feed and permeate channels are given by the conservative two-dimensional forms of the continuity and Navier-Stokes equations and transport equations respectively. Fluid flow within these channels are assumed to be within the laminar regime, the fluid is assumed to be completely incompressible with constant viscosity. Buoyancy forces are assumed to be negligible, which is suitable for short channels flows [5]. Solute flows are assumed to have constant diffusivity within these channels. The system is also assumed to be isothermal.

For flow in narrow empty channels, laminar flows are assumed for Reynolds number, Re, less than 2000 [37]. It is reported in the literature that there is departure from the laminar flow regime at very low Reynolds numbers for spacer filled channels [38-39]. This modified Reynolds number, Re_m , is defined on the basis of superficial velocity and spacer filament diameter. Laminar regime persists for $Re_m < 200$ [38-39]. Manufacturers prescribe that the average velocity and pressure drop for spacer filled channels do not exceed 0.34 ms⁻¹ and 0.6bar/m; this corresponds to Reynolds number, defined on a similar basis, of less than 200 [38]. For all simulations, laminar flows are, therefore, prescribed for both spacer free and spacer filled channels.

2.2 Fluid Flow through the Membrane

The permeate flow, J_w , across the active and porous support layers of the membrane is given by the Eq. (1):

$$J_w = A(1+\alpha) \left(\Delta P - \frac{1}{(1+\alpha)} \Delta \pi \right) \tag{1}$$

Equation 1 is a modified version of the equation for permeate flow derived in solution-diffusion model [2]. The membrane is similarly assumed to be non-porous. A is the pure water permeability, α is a dimensionless semi-empirical relationship that accounts for the effects of pressure, cross-flow velocity of the solute and concentration on solvent flow through the membrane. ΔP is the hydraulic pressure differential across the membrane. It is calculated using the pressure in the feed channel at the membrane surface, $P_{f,m}$, and the pressure in the permeate channel at the membrane surface, $P_{p,m}$. Pressure decreases moving along the membrane surfaces as a result of frictional losses across the channel. As such, ΔP must be calculated for the length of the membrane. It is given by the Eq. (2):

$$\Delta P = P_{f,m}[0 \to L, 0] - P_{p,m}[0 \to L, h1 + h2]$$
(2)

 $\Delta\pi$ is the osmotic pressure differential across the active and porous supports layers of the membrane. It is calculated using the concentration in the feed channel at the membrane surface, $c_{f,m}$, and the concentration in the permeate channel at the membrane surface, $c_{p,m}$. $\Delta\pi$ must be calculated for the length of membrane. It is estimated using Van't Hoff approximation for osmotic pressure of low concentration solutions and is given by Eq. (3):

$$\Delta \pi = \emptyset RT (c_{f,m}[0 \to L, 0] - c_{p,m}[0 \to L, h1 + h2])$$
(3)

Equation 1 includes a semi-empirical relationship that accounts for the effects of inlet pressure, inlet concentration and cross-flow velocity of solutes on solvent across the membrane which is not typically included in the derivation of equations for solvent flows [3, 5, 12-16]. The introduction of the semi-

empirical coefficient within the equation for solvent flow applies only to the hydraulic pressure differential term. Consider the following. The equation used to calibrate pure water permeability coefficient, A, is given by Eq. (4) [31].

$$J_w = A\Delta P \tag{4}$$

In the absence of any concentration, variability of solvent permeation would be as a direct result of inlet pressure. As average cross-flow velocities are a direct result of the pressure differential across the feed channel, this effect is taken into account in Eq. (4) through the calculation of ΔP . The application of the semi-empirical relations is as shown in Eq. (5).

$$J_w = A\Delta P + \alpha A\Delta P \tag{5}$$

Where now a concentration gradient exists, pressure induces the concentration gradient across the membrane [14]. Consideration for the effects of the crossflow velocity and concentration of the solute must be given. They enter into the equation via the semi-empirical relationship. These are seen as "impedances/admittances" towards solvent flow. Eq. (5) is rewritten as shown.

$$J_{w} = A\Delta P + \alpha A\Delta P - A\Delta \pi = A(1+\alpha) \left(\Delta P - \frac{1}{(1+\alpha)}\Delta\pi\right)$$
(6)

Corrections to the osmotic pressure differential are introduced through calculations of the concentration gradient across the membrane.

2.3 Salt Flows through the Membrane

Salt rejection within the membrane is assumed to occur only in the active layer of the membrane. Within the active layer of the membrane, solute transport is assumed to be completely diffusive, consistent. Solute transport is given by Eq. (7):

$$(1+\beta)D_{s(m)}\left(\frac{\partial^2 c_m}{\partial x^2} + \frac{\partial^2 c_m}{\partial y^2}\right) = 0$$
(7)

 β is a dimensionless semi-empirical relation that accounts for the effects of inlet pressure on solute flow through the active layer of the membrane. It is noted that at significantly higher average cross-flow velocities, there may be convection of solute from the polarisation layer [21]. This is not considered in this model. Salt flow transport in the porous support is given by the steady-state conservative form of the convection-diffusion equation, Eq. (8):

$$u\frac{\partial c_{ps}}{\partial x} + v\frac{\partial c_{ps}}{\partial y} = D_{s(ps)}\left(\frac{\partial^2 c_{ps}}{\partial x^2} + \frac{\partial^2 c_{ps}}{\partial y^2}\right)$$
(8)

2.4 Boundary Conditions

2.4.1 Fluid Flow

The pressure and velocity at the feed inlet and outlet are given by the following:

Feed Inlet	Feed out
$p(0,0\to h1)=p_{in}$	$v(L, 0 \to h1) = v_{out}$

At the membrane walls, the velocity is constrained perpendicular to the membrane. Depending on the configuration of the membrane, the velocity components at the wall are given by:

Top of feed channel	Bottom of Feed Channel
$u(0 \to L, y_{h1}) = 0$	$u(0 \to L, 0) = 0$



$$v(0 \to L, y_{h1}) = 0 \qquad \qquad v(0 \to L, 0) = J_w$$

At the porous support layer and permeate channel interface, fluid flow is again constrained perpendicular to the porous support layer. At the permeate outlet, pressure is set to atmospheric.

2.4.2 Solute Flow

The concentration at the inlet is set to:

$$c(0,0 \rightarrow h1) = c_{in}$$

Solute flow is convected at the same rate as the fluid flow within the feed and permeates channels. Within the membrane, solute flow is convected only in the porous layer:

Feed/Permeate Channels	Membrane Channel		
$u_s = u_{f/p}$	$u_s = 0$		
$v_s = v_{f/p}$	$v_s = J_w$		

The solute diffusion coefficients, $D_{s(t/ps/p)}$, in the feed, permeate channels and the porous support layer is set equal to the diffusion coefficient of NaCl in pure water. It is assumed the porous support layer plays no role in the rejection of the solute. At the outlets of the feed and permeate channel, the solute is assumed to be transport out the model by convention, that is the diffusive transport is considered negligible.

3. Simulations and Validation

Simulations were run on a DELL Precision T7500 computer with dual 2.80GHz Xeon processors, 96GB of RAM and Windows 7 operating system. The software used in the simulations was COMSOL Multiphysics versions 5.0 and 5.1. Within the COMSOL environment, the laminar flow node, Darcy's Law node and transport of dilute solute node were used in the simulation.

3.1 Geometry

The geometry used in these simulations is shown in Fig. 1. The dimensions were as reported in [27] and [40]. As the thicknesses of the various membrane layers where unavailable, typical values found in the literature were used. It should be noted that the inlet and outlet of the feed channel were extended to mitigate the development of stress concentrations at the entry and exits. Additionally, as the woven fabric layer of the membrane is assumed to play no role in solvent and solute transport, it is treated as being part of the permeate channel.



Figure 1: Geometry Used in Simulations

3.2 Material Properties

The material properties and dimension used in the simulations are shown in Table 1.

Table 1: Material Properties

Constant Units Value

A ^[27]	m $Pa^{-1}s^{-1}$	2.15×10^{-11}
φRT	J/mol	4988.4
$ ho^{[19]}$	kg m ⁻³	$1.0 \ge 10^3$
$\mu^{[19]}$	Pa s ⁻¹	0.9 x 10 ⁻³
$D_{s (f/p/ps)}^{[41-43]}$	$m^2 s^{-1}$	1.6 x 10 ⁻⁹
$L^{[40]}$	m	1.46 x 10 ⁻¹
$h1^{[44-46]}$	m	5.0 x 10 ⁻⁷
$h2^{[44-46]}$	m	5.0 x 10 ⁻⁵
h3 ^[44-46]	m	1.5 x 10 ⁻⁴
$h4^{[40]}$	m	1.73 x 10 ⁻³

3.3 Mesh Specification

The mesh was manually specified for the feed, membrane and permeate channels. Both triangular and quadrilateral elements were used in the mesh. Areas critical to the reverse osmosis process were meshed with a very fine mesh. These included the individual membrane layers and regions just above the membrane. Within the feed channel, an internal boundary was set at a distance of 5% of the feed channel height, h1, from the feed channel and membrane interface. This was to ensure that a sufficiently small mesh was defined to accurately measure the effects of the concentration polarisation layer [23, 47], while optimising the number of elements for a grid independent mesh.

3.4 Grid Independence

Grid impendence was determined based on two criteria: (1) the convergence of outlet parameters to less than 1% of successive mesh refinements for constant inlet conditions; (2) an increase in the minimum element quality as given within the COMSOL environment. Grid independence was found on a mesh of approximately 6.7 million elements. Fig. 2 shows the convergence of average permeate flow and average permeate concentration for increasing mesh density.

3.5 Validation

Simulations were run for the operating conditions as described in Kim & Hoek [26], who considered three inlet operating pressure (790 kPa, 1136 kPa and 1481kPa), for three average feed channel velocities (0.017 m s⁻¹, 0.042 m s⁻¹ and 0.068 m s⁻¹) and three inlet concentrations (10 mol m⁻³, 20 mol m⁻³ and 50 mol m⁻³). For these simulations, $D_{s(m)}$ was set to 380.3 x 10⁻¹⁵ m² s⁻¹ having been calibrated for one of the experimental data points, α and β ranged between -0.176 to 0.038 and -0.197 to 0.072 respectively.

Fig. 3 - 5 compare the simulations results for average permeate flows and permeate concentration versus experimental average permeate flows and permeate concentration. The simulations agreed with the experimental data points. The average predicted error for permeate flow across all simulations was -0.6% \pm 2.6%. The average predicted error for permeate concentration was 0.7% \pm 7.6%.

4. Feed Spacers

Feed spacers are used to keep the membrane(s) apart from the flow channel. It is therefore practical to assume that the channel height is twice the feed spacer diameter. Similar consideration was given by Amokrane, et al. [48] and Wardeh & Morvan [33]. In these simulations, the filaments were assumed to be circular with a diameter (D) of 0.865 mm. These filaments were assumed to be completely rigid. The distance between the spacers was determined using the geometrical ratio (G.R.) as defined by [29] and is given by Eq. (9).

$$G.R. = \frac{L_f}{D} \tag{9}$$

In the above equation, L_f is the distance between parallel filaments and D is the diameter of the filament. It is further noted that the G.R. is typically 7-9 [29]. Using these G.R. values, L_f can be estimated and the number of filaments determined. For the length of channel under consideration and G.R. of 7, 8 and 9 correspond to 24, 21 and 19 filaments within the feed channel. The angle between crossing filaments and the flow attack angle were assumed to be 90°. Three (3) spacer configurations were considered: submerged, cavity and zigzag.



Figure 2: Grid Independence. Boundary and Simulation Conditions: $p_{in} = 1136.3$ kPa, $u_{avg} = 0.042$ m s⁻¹, $c_{in} = 50$ mol m⁻³, $D_{s(m)} = 400.0$ x 10^{-15} m² s⁻¹, α and $\beta = 0$





Figure 3: Comparison of Simulated Average Permeate Flows and Permeate Concentration versus Experimental Permeate Flows and Permeate Concentration: Avg. Crossflow Velocity of 0.017 m s⁻¹



Figure 4: Comparison of Simulated Average Permeate Flows and Permeate Concentration versus Experimental Permeate Flows and Permeate Concentration: Avg. Crossflow Velocity of 0.042 m s⁻¹





Figure 5: Comparison of Simulated Average Permeate Flows and Permeate Concentration versus Experimental Permeate Flows and Permeate Concentration: Avg. Crossflow Velocity of 0.068 m s⁻¹

The average permeate flow, average permeate concentration, average feed channel velocity, feed channel pressure losses and maximum wall shear stress for the three spacer types are shown in Table 2. Independent of spacer type, there is an increase in permeate flows across the membrane and a decrease in the average permeate concentration as compared to spacer free channels. The pressure loss across the feed channel increase with the presence of spacers; larger losses seen for submerged types. The submerged spacer type causes a larger pressure loss; cavity and zigzag spacer types have similar losses. Similar pressure pattern losses are reported in the literature [33-34]. Larger wall shear stress is seen for the submerged type spacer than for either cavity or zigzag types. Similar observations are found in the literature [33, 49].

Spacer Type	Geometric Ratio	Sim J _w (µm/s)	Sim. C _p (mol/m ³)	Avg. Channel Velocity (m/s)	Feed Channel Pressure Loss (Pa/m)	Max. Wall Shear Stress (N m ⁻²)
No Spacer	n/a	11.97	3.738	0.0420	158.5	0.1977
Submerged	7	13.15	3.064	0.0453	1039.0	1.1781
	8	13.03	3.128	0.0448	930.3	1.1777
	9	12.92	3.187	0.0446	857.1	1.1773
	7	13.33	3.688	0.0426	488.7	0.2197
Cavity	8	13.67	3.481	0.0422	446.6	0.2432
	9	13.83	3.367	0.0419	418.0	0.2435
Zigzag	7	14.30	3.015	0.0427	477.8	0.5523
	8	14.25	3.005	0.0422	440.6	0.5549
	9	14.19	2.983	0.0419	415.5	0.5585

Table 2: Simulation Results for Submerged, Cavity and Zigzag Spacer Types.

Note: Boundary and Simulation Conditions: $p_{in} = 1136.3$ kPa, $u_{out} = 0.04165$ m s⁻¹, $c_{in} = 50$ mol m⁻³, $D_{s(al)} = 380.3$ x 10^{-15} m² s⁻¹, $\alpha = -0.0266$ and $\beta = 0.0214$.

5. Conclusion

In this paper, the limitations restricting the use of currently available models for the specific application of wave powered desalination were considered. Fully coupled mass-momentum equations were used to describe the hydrodynamics with the feed and permeate channels. Equations were specified to describe for solvent and solute flows across the individual layers comprising thin-film membranes. Additionally, semi-empirical relationships were developed to account for the effects of inlet pressure, cross-flow velocity and concentration on solvent and solute flows across the membrane. The model was validated against published experimental data. The predicted errors for simulations of solvent and solute flows were found to be 0.6% \pm 2.6% and 0.7% \pm 7.6% respectively. The effects of three feed spacer types, submerged, cavity and zigzag, on solvent and solute flows were then considered. Key metrics, such as the feed channel pressure loss and maximum wall shear stress at the membrane surface, agreed with the literature.

Acknowledgements

The authors wish to acknowledge the financial grant provided by BP of Trinidad and Tobago for the conduct of this research.

Nomenclature

A: Pure Water Permeability Coefficient $(m^2 s kg^{-1})$ P: PressB: Salt Permeability Coefficient (m/s)R: Universec: Solute Concentration $(mol m^{-3})$ Re: ReyrJw: Permeate Flow through the Membrane $(m s^{-1})$ Rem: Modelh1: Feed Channel Height (m)T: Temph2: Active Layer of Membrane Thickness (m)u: Transh3: Porous Support Layer of Membrane Thickness (m)v: Axialk: Effective Intrinsic Permeability (m^2) x: axial ofL: Length of the membrane (m)y: normation

m: Active Layer of the Membranef: Feed Channelps: Porous Support of the Membrane

Greek Letters

- α: Semi-Empirical Relationship Affecting Water Flow across Membrane
- β: Semi-Empirical Relationship Affecting Solute Flow across Membrane
- μ: Absolute Viscosity (Pa s⁻¹)
- ρ: Density (kg m⁻³)
- υ: Kinematic Viscosity (m² s)
- Ø: van't Hoff index

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P: Pressure at Location x, y (Pa) R: Universal Gas Constant (J/mol·K) Re: Reynolds Number Re_m: Modified Reynolds Number T: Temperature (K) u: Transverse Velocity Component (m s⁻¹) u_{avg}: Average Feed Channel Velocity (m s⁻¹) v: Axial Velocity Component (m s⁻¹) x: axial coordinate (m) y: normal coordinate (m)

p: Permeate Channel b: Bulk Solution

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